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# The Effect of Coulomb Drag on the Mean Motion of Low Earth Orbit Satellites

## Classical Gravitation at Two Body Problems

The theoretical results all stem from the classical gravitational two body problem. The two bodies are clearly the earth and satellite [1].

$$GM = a(t)^2 n(t)^3 \quad (1)$$

Then, taking the time derivative can be taken and by doing some rearranging of the equation the result becomes [1].

$$\frac{dn}{dt} = -\frac{3n(t)}{2a(t)} \frac{da}{dt} \quad (2)$$

This is a widely recognized result seen in Haranas 2010 [1]. This equation allows us to represent the change in mean motion with respect to time as related to the change in semi major axis in time by a time dependent function. This, can then by a simple substitution of Gauss' relations which are as follows. This substitution will later yield the theoretical model for coulombic drag of low earth orbit satellites [2].

$$\frac{da}{dt} = \frac{n(t)a(t)^2}{(1-e^2)} C_{D,e} \frac{(\pi R^2)}{m_s} m_e N (1+e^2 + 2e \cos(f))^{\frac{3}{2}} \quad (3)$$

## The Predicted Effects of Coulomb Drag on the Mean Motion of Low Earth Orbit Satellites

From the previous section we know important equations which are useful for many parameters of objects in orbit. By a simple substitution of equation 3) into equation 1) yields a differential equation which will become the main predictive equation. However, it is also necessary to do a transform of the fundamental anomaly with the eccentric anomaly. This transformation is done with the following transformation [1].

$$\cos(f) = \frac{\cos(E) - e}{1 - e \cos(E)} \quad (4)$$

This will allow us to express the change of mean motion where the variable parameter is the eccentric anomaly. The eccentric anomaly will be a parameter in radians and for one full revolution will move from 0 to  $2\pi$ , or a transformation can be done which will change it from 0 to 360 degrees. By doing this substitution we will yield a rather large and cumbersome equation which will have the eccentricity (e) expressed as a cubic equation where we can omit the third order term, this is because the eccentricity is a value less than 1 so the cubed is much less than 1. This now will make our result more simplified while sacrificing negligible accuracy. However, this form will be a differential with respect to time, where it is more valuable to have it as a variable with respect to the eccentric anomaly. In order to do this it is necessary to another transformation which will be the following [2].

$$\frac{dE}{dt} = \frac{na}{r} = \frac{n}{1 - e \cos(E)} \quad (5)$$

The result after this substitution will be a linear separable differential equation which will have the limits of integration between 0 and  $2\pi$ , which has the following form.

$$\int \frac{dn}{n} = \frac{-3aC_{D,e}Nm_e R^2 \pi}{2m_s \sqrt{1-e^2}} \int_0^{2\pi} \left(1 + \frac{e}{1-e^2} \cos(E)\right) dE \quad (6)$$

By solving this relatively simple equation we yield the following result, which is our theoretical result.

$$\frac{n(E)}{n_0} = e^{-\frac{3aC_{D,e}Nm_e(\pi R)^2}{m_s \sqrt{1-e^2}}} \quad (7)$$

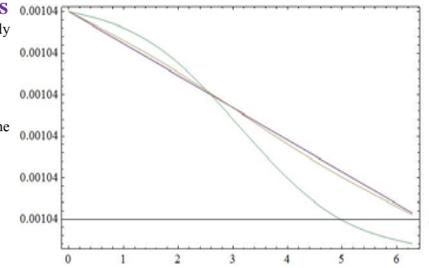


Figure 1. The change of mean motion LEO satellites for one full revolution at 250km in altitude.

## Numerical Results

We will take the values of eccentricity of 0, 0.01, 0.13, and 0.5. A 0 eccentricity will represent the orbit of a satellite which is circular, and 0.01 is nearly circular. The eccentricity of 0.13 represents a very usual eccentricity of typical low earth orbit satellites. The eccentricity of highly elliptical satellites. We will also use parameters of radius of the satellite as 0.5 meters which is a very typical size of many modern day satellites. However, a satellite such as Hubble which is the size of a school bus is much bigger, but this is not the case we are considering which are usually smaller communication satellites. We will consider several values for the altitude which are 250km, 500km, 1000km, and 20,000km in altitude [2]. This corresponds to a semi major axis of 6671km, 7171km, 7671km, and 26621km respectively. Other parameters which are used in these calculations are the mass of the satellite which is denoted by  $m_s$  this will be 1000 kg which is average size of most low earth orbit satellites. Then is also the mass of the electron  $m_e$  which is  $9.11 \times 10^{-31} kg$  [5]. The last of the values which are relevant to the numeric calculations is the density of electrons in the range of semi major axis' which we are finding solutions for, this value is denoted by  $N$  we know the value is of the order of  $10^{12} m^{-3}$  [2].

Table 1 Values for the Drag Coefficient with Respect to Altitude [2]

Height	250km	500km	1000km	20,000km
Height (km)	250	500	800	$2 \times 10^4$
$C_{D,e}$	$7 \times 10^{-5}$	0.32	6.1	$7.7 \times 10^4$
$C_{D,e}$	$2.6 \times 10^4$	0.5	1.38	$1.1 \times 10^4$

Table 2 The Numeric results for mean motion per revolution for eccentricity  $e = 0.13$  at various orbital altitudes.

Height	250km	500km	1000km	20,000km
$a_0$ (km)	6621	6871	7171	26371
$n(E)/n_0$	0.99921	0.9999999	1.0	1.0



## Conclusions

To conclude, by using the well known gravitational two body problem and Gauss' Planetary equations it is possible to predict the orbit of a low orbit earth orbit satellite, and its effect on its semimajor axis. The work done in this paper predicted a simple first order linear differential equation for the mean motion of the low orbit satellite, that resulted to a decaying exponential solution. Larger satellites may play a role, in how Coulomb drag effects the mean motion of satellites. However, the eccentricity of the orbit and the orbital altitude are the orbital elements that significantly affect affecting the Coulomb drag.

## References

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